

The two expressions may be combined into

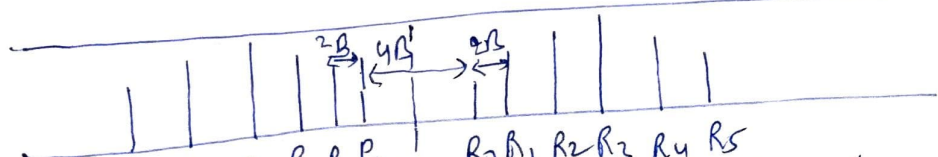
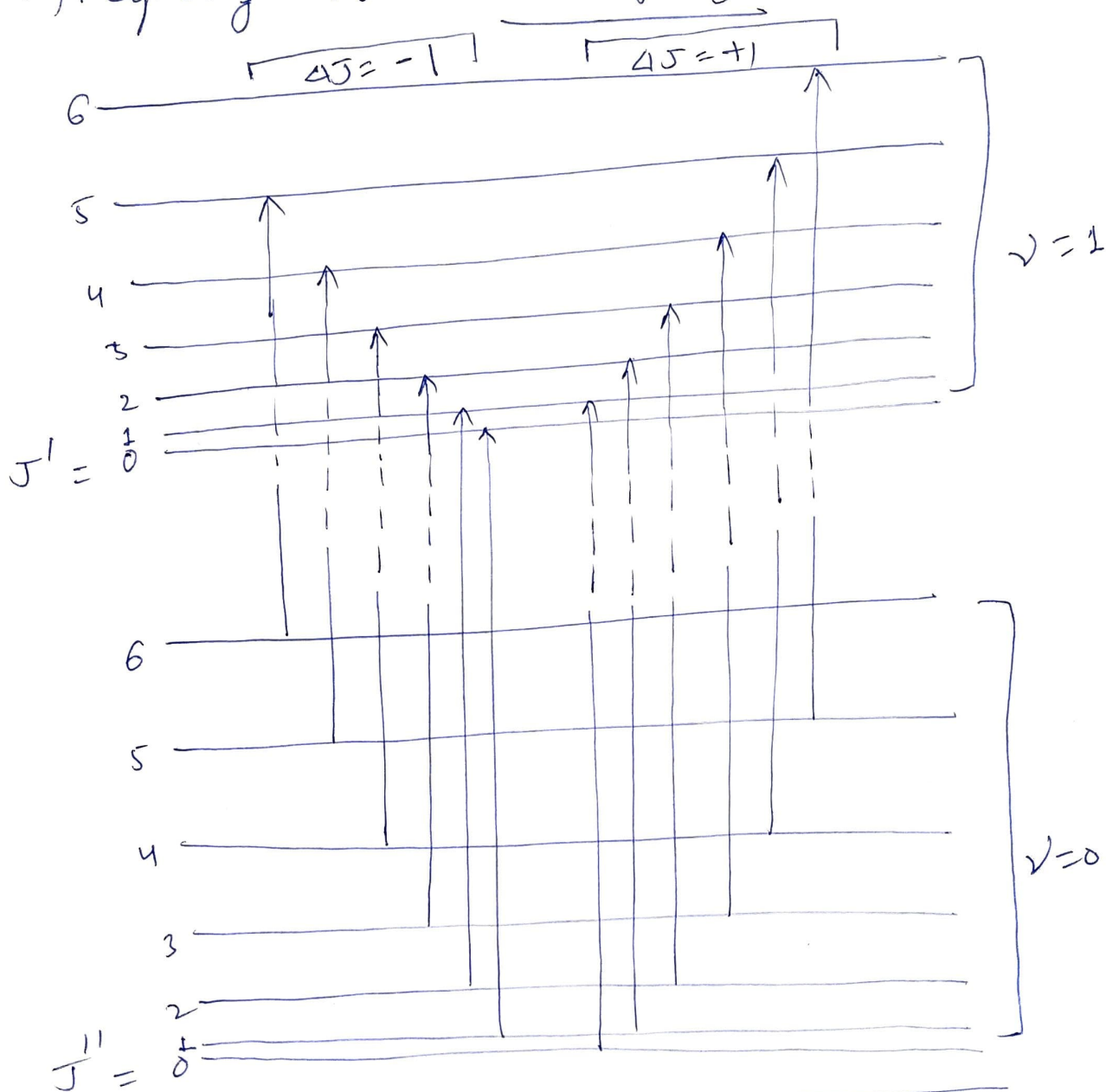
$$\Delta E_{J',v} = \bar{\nu}_{\text{spect}} = \bar{\omega}_0 + 2Bm \text{ cm}^{-1} \quad \text{--- 20(c)}$$

$$m = \pm 1, \pm 2 \dots$$

$m \rightarrow$ can not be zero, because that implies J' or J'' to be -1 .

$m \rightarrow$ replacing $J''+1$ in eqⁿ 20(a) and $J'+1$ in 20(b) has positive values for $\Delta J = +1$ and is negative for $\Delta J = -1$

The frequency $\bar{\omega}_0 \rightarrow$ band ~~of~~ origin and band centre.



Transitions between Rotational - vibrational energy levels ~~exists~~ of a diatomic molecule together with spectrum arising

Eqⁿ 20(c) → represents combined vibration-rotational Spectrum.

Consists of equally spaced lines (spacing = 2B) on each side of the band origin $\bar{\omega}_0$, since $m \neq 0$, the line $\bar{\omega}_0$ itself will not appear.

Lines to the low frequency side of $\bar{\omega}_0$ → negative m (i.e. $\Delta J = -1$) → P Branch

High-frequency side (m positive, $\Delta J = +1$) → R Branch

Lines arising from	$\Delta J = -2$	-1	0	$+1$	$+2$
Called:		O	P	Q	R S
					branches

The inclusion of the centrifugal distortion constant D leads to the following expressions for the spectrum

$$\Delta E = \bar{\nu}_{\text{spect}} = \bar{\omega}_0 + 2Bm - 4Dm^3 \text{ cm}^{-1} \quad \text{--- (21)}$$

$(m = \pm 1, \pm 2, \pm 3, \dots)$

$B \sim B_0$ of the order 10 cm^{-1} or less

$D \sim$ only some 0.01 of B

Good infra-red spectrometer has a resolving power $\approx 0.5 \text{ cm}^{-1}$, $D \rightarrow$ negligible to a very high degree of accuracy.

Anharmonicity factor → not negligible it affects the position of the band origin ($\bar{\omega}_0 = \bar{\omega}'_e(1 - 2x_e)$)

and the selection rules to include $\Delta v = \pm 2, \pm 3$ etc.

Also allows the appearance of overtone bands having identical rotational structure.

Break down of Born-Oppenheimer Approximations:

The Interaction of Rotations and Vibrations

So far we have assumed \rightarrow Vibrations and rotations can proceed quite independently of each other.

A molecule vibrates some 10^3 times during the course of a single rotation \rightarrow the bond length and hence the moment of inertia and B constant also changes continually during the rotation.

If vibration is simple harmonic \rightarrow mean bond length will be same as equilibrium bond length and will not vary with vibrational energy.

The rotational constant $B \rightarrow$ depends as $\frac{1}{r^2}$, the average value of $\frac{1}{r^2}$ is not same as $\frac{1}{r_{eq}^2}$.

Increase in vibrational energy \rightarrow Increase in the vibrational amplitude \rightarrow hence value of B will depend on the v quantum number.

Case of anharmonic vibrations \rightarrow more complex.

Increase in vibrational energy \rightarrow increase in the average bond length. The rotational constant then varies even more with vibrational energy.

In general since r_{av} increases with vibrational energy \rightarrow B is smaller in upper vibrational state than in the lower.

In eqⁿ form

$$B_v = B_e - \alpha \left(v + \frac{1}{2} \right) \quad (22)$$

$B_v \rightarrow$ rotational constant in vibrational level v . (42)

$\alpha \rightarrow$ small positive constant for each molecule

Fundamental vibrational change, i.e. change

$v=0$ to $v=1$ $v=0 \rightarrow v=1$

$$B_0 > B_1$$

$$\Delta E = \sum_{J', v=1} - \sum_{J'', v=0}$$

$$= \bar{\omega}_0 + B_1 J'(J'+1) - B_0 J''(J''+1) \text{ cm}^{-1}$$

$$\text{where } \bar{\omega}_0 = \bar{\omega}_e(1-2\alpha_e).$$

1. $\Delta J = +1$ $J' = J'' + 1$

$$\Delta E = \bar{\nu}_R = \bar{\omega}_0 + (B_1 + B_0)(J'' + 1) + (B_1 - B_0)(J'' + 1)^2 \text{ cm}^{-1} \quad \text{--- 23(a)}$$

$(J'' = 0, 1, 2, \dots)$

and

2. $\Delta J = -1$ $J'' = J' + 1$

$$\Delta E = \bar{\nu}_P = \bar{\omega}_0 - (B_1 + B_0)(J' + 1) + (B_1 - B_0)(J' + 1)^2 \text{ cm}^{-1} \quad \text{--- 23(b)}$$

$(J' = 0, 1, 2, \dots)$

$\bar{\nu}_P$ and $\bar{\nu}_R \rightarrow$ wavenumbers of P and R branch lines respectively.

So $\bar{\nu}_{P,R} = \bar{\omega}_0 + (B_1 + B_0)m + (B_1 - B_0)m^2 \text{ cm}^{-1}$
 $(m = \pm 1, \pm 2, \dots)$

Positive m value \rightarrow R branch 23(c)

Negative m value \rightarrow P branch

Ignoring vibration-rotation interactions, $B_1 = B_0$,
eqⁿ 23(c) simplifies to 20(c).

(43)

Since $B_1 < B_0$ the last term in eqⁿ 23 (c) is always negative \rightarrow irrespective of the sign of m ,
 \rightarrow the effect on the spectrum of a diatomic molecule
 \rightarrow crowd the rotational lines more closely together with increasing m on the R branch side and the P branch side lines become more widely spaced as (negative) m increases.

Normally B_1 and B_0 differ slightly \rightarrow effect is marked only for high m values.

The Vibrations of Polyatomic Molecules

- * The number of fundamental vibrations and their symmetry
- * The possibility of overtone and combination bands
- * The influence of rotation on the spectra

Fundamental Vibrations and their Symmetry

Consider a molecule containing N atoms:
We can refer to the positions of each atom by specifying three coordinates (i.e., x , y and z Cartesian coordinates). Total no. of coordinate values is $3N \rightarrow$ molecule has $3N$ degree of freedom \rightarrow since each coordinate value may be specified quite independently of the others. Once all $3N$ coordinates have been fixed \rightarrow the bond distances and bond angles of the molecule are also fixed \rightarrow no further arbitrary specifications.